

This document is posted with the permission of the author Peter May. It was hand-written by Peter, most likely in 1963 or 1964. It describes the cohomology of $A(2)$, which later came to be known as the Adams E_2 -page for tmf . The computation is correct in all respects (as far as I have noticed). It gives a complete roadmap for the May spectral sequence, including all differentials and all hidden multiplicative structure. Although it was never published, this document slightly precedes the work of Shimada and Iwai that was published in 1967 (and received by the journal in 1966).

This document represents a very early computational hint of the existence of a spectrum that would be constructed much later, and even later be named topological modular forms.

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$p=2$ $A_n \subset A$ generated by $\{4^1, \dots, 4^{2^n}\}$ ($= A(n+1)$ in my paper, chapter I)

$E^\circ A_n = VL_n$ where L_n has basis $\{P_j^i \mid i+j \leq n+1\}$, $\bar{X}(L_n)^* = P \{R_j^i \mid i+j \leq n+1\}$ with

$dR_j^i = c_{ij} = \sum R_k^i R_{j+k}^{i+k}$. Higher differentials computable by pull-back from $E_r A$.

(1) $A_0 = E\{4^1\}$ hence $H^*(A_0) = P\{h_0\}$

(2) $H^*(E^\circ A_1) = P\{h_0, h_1, b_2^0\} / (h_0 h_1)$ and $d_2(h_1^2) = h_1^3$ since $b_2^0 = \langle h_1, h_0, h_1, h_0 \rangle$. Then

$E_\infty = E_3 = P\{h_0, h_1, X, w\} / (h_0 h_1, h_1^3, h_1 X, X^2 + h_0^2 w) = H^*(A_1)$, where we have

$h_0 \in (1,1)$, $h_1 \in (1,2)$, $X = h_0 b_2^0 = \langle h_0, h_1, h_1^2 \rangle \in (3,7)$, $w = (b_2^0)^2 = \langle h_1^2, h_1, h_1^2, h_1 \rangle \in (4,12)$.

Thus $H^*(A_1) = P\{w\} \otimes B\{h_0^i, h_0^i X, h_1, h_1^2 \mid i \geq 0\}$ as a left $P\{w\}$ -module.

(3) $H^*(E^\circ A_2) = P\{h_0, h_1, h_2, h_0(1), b_2^0, b_2^1, b_3^0\} / I$ where I is the ideal generated by

$h_0 h_1, h_1 h_2, h_0 h_0(1) + h_2 b_2^0, h_2 h_0(1) + h_0 b_2^1, h_0(1)^2 + h_1^2 b_3^0 + b_2^1 b_2^0$ (hence $h_0^2 b_2^1 + h_2^2 b_2^0 \in I$).

$$h_0(1) \leftrightarrow R_2^0 R_2^1 + R_1^1 R_3^0$$

Thus $E_2 A_2$ admits the following Z_2 -basis ($a > 0, l \geq 0, b \geq 0, c \geq 0, k \geq 0, i \geq 0, j \geq 0, e = 0 \text{ or } 1$)

- (i) $(b_2^1)^b (b_3^0)^c h_0^i h_2^j, i+j > 0$
- (ii) $(b_2^1)^b (b_3^0)^c h_0^{2l(1)} h_0^i h_2^e, i+e > 0, l > 0$
- (iii) $(b_2^0)^a (b_3^0)^c h_0^{2l(1)} h_0^i h_2^e, i+e > 0, a > 0$
- (iv) $(b_2^1)^b (b_3^0)^c h_1^k h_0^l(1)$
- (v) $(b_2^0)^a (b_3^0)^c h_1^k h_0^l(1), a > 0$

$d_2(b_2^0) = h_1^3 + h_0^2 h_2, d_2(b_2^1) = h_2^3, d_2(b_3^0) = h_1 b_2^1$, and $d_2 h_0(1) = h_0 h_2^2$. Let

$w = (b_2^0)^2, \bar{p} = (b_3^0)^2, a = h_1 h_0(1), b = h_0^2(1), c = b_2^1 h_0(1), d = (b_2^1)^2, \alpha = h_0 b_3^0, \nu = h_2 b_3^0$;

writing out d_2 in terms of our basis, we find that E_3 admits as basis:

- (a) $P\{w, d, \bar{p}\} \otimes E\{b\} \otimes B\{1, h_0, h_1, h_2, h_0 h_2, h_2^2, c, h_1 c, \alpha, \nu, h_0 \nu\}$
- (b) $P\{w, d, \bar{p}\} \otimes B\{h_0 \alpha, h_2 \nu\}$
- (c) $P\{w, \bar{p}\} \otimes B\{h_1^2, h_1^3, a, h_1 a, h_0^i, h_0^i \alpha \mid i \geq 2\}$

$d_4(\bar{p}) = h_2 d$ and using the relations in E_3 , we find that $E_5 = E_{\infty}$ admits as basis:

$$(d) P\{\omega, d, \beta\} \otimes E\{b\} \otimes B\{1, c, \alpha, v, \alpha v, v^2, \gamma, \gamma c\}$$

$$(e) P\{\omega, \beta\} \otimes B\{h_0^{(1)}, h_0^{(2)}, h_0^{(3)}, h_0^{(4)}, h_1, h_1^2, h_2, h_0 h_2, h_0^2 h_2, h_2^2, a, h_1 a, h_0 b, h_0 v, h_0^2 v, h_0 c, h_0^2 c, h_2 v, h_0 h_2 v, h_0 d, h_0^2 d, h_1 d, h_0 v b, h_0 \alpha d, h_0^2 \alpha d, h_0 b d, h_1 \gamma, h_1^2 \gamma, \delta, h_1 \delta \mid i=0\}$$

where $\beta = \bar{\beta}^2 = (b_3^0)^4$, $\gamma = h_1 \bar{\beta} = h_1 (b_3^0)^2$, and $\delta = a \bar{\beta} = h_1 h_0 (1) (b_3^0)^2$. Thus E_∞ is generated as an algebra by $\{\omega, d, \beta, h_0, h_1, h_2, a, b, c, \alpha, v, \gamma, \delta\}$ where

$h_i \in (1, 2^i)$, $i = 0, 1$, and 2 , and

$$\omega = \left\langle (h_1^2, h_0), \begin{pmatrix} h_1 & h_0 \\ h_0 h_2 & h_1^2 \end{pmatrix}, \begin{pmatrix} h_1^2 & h_0 \\ h_0 h_2 & h_1 \end{pmatrix}, \begin{pmatrix} h_1 \\ h_0 h_2 \end{pmatrix} \right\rangle \in (4, 12)$$

$$d = \langle h_2, h_2^2, h_2, h_2^2 \rangle = \langle h_2, h_2^2, h_2^2, h_2 \rangle \in (4, 24)$$

$$\beta = \langle h_2, d, h_2, d \rangle \in (8, 56)$$

$$a = \langle h_1, h_0, h_2^2 \rangle \in (3, 11)$$

$$b = \langle h_0, h_2^2, h_0, h_2^2 \rangle = \langle h_0, h_2^2, h_2^2, h_0 \rangle \in (4, 18)$$

$$c = \langle h_0, h_2^2, h_2, h_2^2 \rangle = \langle h_0, h_2^2, h_2^2, h_2 \rangle \in (4, 21)$$

$$\alpha = \langle h_0, h_1, h_2, h_2^2 \rangle \in (3, 15)$$

$$v = \langle h_2, h_1, h_2, h_2^2 \rangle \in (3, 18)$$

$$\gamma = \langle h_1, h_2, d \rangle \in (5, 30)$$

$$\delta = \langle a, d, h_2 \rangle \in \langle a, h_2, d \rangle \text{ (}\alpha d \text{ intertwining)} \in (7, 39)$$

a defining system of relations for E_∞ is the following:

$$h_0 h_1 = 0, h_1 h_2 = 0, h_0^2 h_2 = h_1^3, h_0 h_2^2 = 0, h_2^3 = 0;$$

$$h_0 a = 0, h_1^2 a = 0, h_2 a = 0, a^2 = 0, ab = 0, ac = 0, ad = 0, \alpha a = 0, va = 0;$$

$$h_0^2 b = h_2^2 \omega, h_1 b = h_0^2 v, h_2 b = h_0 c, h_0 \alpha b = h_2 v \omega, \alpha^2 b = v^2 \omega, vb = \alpha c, b^2 = \omega d, bd = c^2;$$

$$h_1 c = h_2^2 \alpha, h_0 d = h_2 c, h_1 d = h_2^2 v, h_2 d = 0, \alpha d = v c;$$

$$h_1 \alpha = 0, h_1 v = 0, h_2 \alpha = h_0 v, h_2 \alpha^2 = h_1^2 \gamma, h_0 v^2 = 0, h_2 v^2 = 0;$$

$$h_0 \gamma = 0, h_2 \gamma = 0, \gamma a = h_1 \delta, \gamma b = \alpha^2 v, \gamma c = \alpha v^2, \gamma d = v^3, \alpha \gamma = 0, v \gamma = 0, \gamma^2 = h_1^2 \beta;$$

$$h_0 \delta = 0, h_1^2 \delta = 0, h_2 \delta = 0, \delta a = 0, \delta b = 0, \delta c = 0, \delta d = 0, \alpha \delta = 0, v \delta = 0, \gamma \delta = h_1 \alpha \beta, \delta^2 = 0, \alpha^4 = h_0^4 \beta$$

$H^*(A_2)$ has generators $\{w, \alpha, v, \rho, \gamma, \delta, a, b, c, d, h_0, h_1, h_2\}$ given by the above Massey products. Clearly $\langle h_0, h_1, h_0 \rangle = h_1^2$, $\langle h_1, h_0, h_1 \rangle = h_0 h_2$, and $\langle h_1, h_2, h_1 \rangle = h_2^2$. Thus

$$h_1 \langle h_0, h_1, \alpha \rangle = \langle h_1, h_0, h_1 \rangle \alpha = h_0 h_2 \alpha = h_1 b \Rightarrow b = \langle h_0, h_1, \alpha \rangle$$

$$h_1 \langle h_2, h_1, \alpha \rangle = \langle h_1, h_2, h_1 \rangle \alpha = h_2^2 \alpha = h_1 c \Rightarrow c = \langle h_2, h_1, \alpha \rangle$$

$$h_1 \langle h_0, h_1, v \rangle = \langle h_1, h_0, h_1 \rangle v = h_0 h_2 v = h_1 c \Rightarrow c = \langle h_0, h_1, v \rangle$$

$$h_1 \langle h_2, h_1, v \rangle = \langle h_1, h_2, h_1 \rangle v = h_2^2 v = h_1 d \Rightarrow d = \langle h_2, h_1, v \rangle$$

We use this to check all dimensionally possible non-trivial extensions.

$$\alpha c = \alpha \langle h_1, h_0, h_2^2 \rangle = \langle \alpha, h_1, h_0 \rangle h_2^2 = b h_2^2 = h_0^2 d$$

$$\alpha \gamma = \alpha \langle h_1, h_2, d \rangle = \langle \alpha, h_1, h_2 \rangle d = c d$$

$$\alpha_1^4 \gamma = \alpha_1^3 c d = \gamma d^2 w \neq 0 \text{ but } h_0^4 \gamma \rho = 0 \Rightarrow \alpha_1^4 = h_0^4 \rho + q^2 w$$

$$v \gamma = v \langle h_1, h_2, d \rangle = \langle v, h_1, h_2 \rangle d = d^2$$

$$v \gamma^2 = \gamma d^2 \neq 0 \text{ but } v h_1^2 \rho = 0 \Rightarrow \gamma^2 = h_1^2 \rho + v^2 d$$

Since $d_2(b_2^0 h_0(1)) = h_1^3 h_0(1)$ and $\{h_2 b_2^0 h_0(1)\} = \{h_0 b\}$, we have $\langle h_1 a, h_1, h_2 \rangle = h_0 b$, hence

$$h_1^2 \delta = h_1 a \gamma = h_1 a \langle h_1, h_2, d \rangle = \langle h_1 a, h_1, h_2 \rangle d = h_0 b d. \text{ Recall that } \delta = \langle h_2, d, a \rangle;$$

since $h_0 a = 0$ in E_2 and $d_2(b_3^0 c) = a d$, $\langle d, a, h_0 \rangle = \alpha c = v b$, and therefore

$$h_0 \delta = \langle h_2, d, a \rangle h_0 = h_2 \langle d, a, h_0 \rangle = h_2 v b = h_0 \alpha d.$$

$$\gamma \delta d = 0 = h_1 a \rho d \text{ but } c d^3 \neq 0, \text{ hence } \gamma \delta = h_1 a \rho;$$

$$\alpha \delta = \langle h_0, h_1, h_2, h_2^2 \rangle \delta \subset \langle h_0, h_1, h_2, 0 \rangle = 0 \text{ (strictly defined, zero indeterminacy)}$$

$$v \delta = \langle h_2, d, a \rangle v = h_2 \langle d, a, v \rangle = 0$$

$$b \delta = h_2 \langle d, a, b \rangle = 0$$

$$c \delta = h_2 \langle d, a, c \rangle = 0$$

$$d \delta = h_2 \langle d, a, d \rangle = 0$$

$$\delta^2 = h_2 \langle d, a, \delta \rangle = 0$$

Since $w(h_0 a) = 6 < 8 = w(w)$, $h_0 a = 0$; since $w(\alpha d) = w(v c) = 15 < 17 = w(\delta)$,

$\alpha d = v c$. There are no further possible non-trivial extensions.

Thus $H^*(A_2)$ admits the following defining set of relations:

$$(i) \quad h_0 h_1 = 0, h_1 h_2 = 0, h_0^2 h_2 = h_1^3, h_0 h_2^2 = 0, h_2^3 = 0$$

$$(ii) \quad h_0 a = 0, h_1^2 a = 0, h_2 a = 0, a^2 = 0, ab = 0, ac = 0, ad = 0, a\alpha = h_0^2 d, av = 0$$

$$(iii) \quad h_0^2 b = h_2^2 w, h_1 b = h_0^2 v, h_2 b = h_0 c, b^2 = wd, bd = c^2, h_0 \alpha b = h_2 vw, \alpha^2 b = v^2 w, vb = \alpha c;$$

$$(iv) \quad h_1 c = h_2^2 \alpha, h_2 c = h_0 d, vc = \alpha d, h_1 d = h_2^2 v, h_2 d = 0, cd = \alpha \gamma, d^2 = v\gamma;$$

$$(v) \quad h_1 \alpha = 0, h_1 v = 0, h_2 \alpha = h_0 v, h_0 v^2 = 0, h_2 v^2 = 0, \alpha^4 = h_0^4 \rho + d^2 w$$

$$(vi) \quad h_0 \gamma = 0, h_1^2 \gamma = h_2 \alpha^2, h_2 \gamma = 0, a\gamma = h_1 \delta, b\gamma = \alpha^2 v, c\gamma = \alpha v^2, d\gamma = v^3, \gamma^2 = h_1^3 \rho + v^2 g;$$

$$(vii) \quad h_0 \delta = h_0 \alpha d, h_1^2 \delta = h_0 b d, h_2 \delta = 0, a\delta = 0, b\delta = 0, c\delta = 0, d\delta = 0, \alpha\delta = 0, v\delta = 0, \gamma\delta = h_1 a \rho, \delta^2 = 0$$